

Investigations on a Universal Relationship Between Optical Emission and Absorption of Complex Molecules in Liquid Solutions

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Based on a universal relationship between the extinction coefficient and the fluorescence intensity in their overlapping region, local temperatures T^* higher than the ambient T were determined for short-lived luminescent molecules of lifetimes from 7 ps to 77 ps. The reason for such a local temperature T^* , which holds also during the emission process, is the non-establishment of statistical equilibrium over the vibrational levels of excited molecules. It is found that the intensity distribution in the fluorescence band depends slightly on the wavelength of the exciting light, which evidences the lack of thermal equilibrium with the vicinal surrounding.

Key words: Universal Relationship between the Absorption and Fluorescence Spectra; Local Temperature; Statistical Equilibrium.

1. Introduction

The universal relation between the extinction coefficient and the fluorescence intensity is based on very general physical assumptions and is one of the most important photoluminescence regularities. As early as in the twentieths Kennard [1, 2] and then Stepanov [3, 4] dealt with that problem. The relation obtained by them can be presented in the form

$$F(\tilde{\nu}) \equiv \ln \frac{f(\tilde{\nu})}{\tilde{\nu}^3 \cdot \varepsilon(\tilde{\nu})} = -\frac{hc\tilde{\nu}}{kT} + \text{const}, \quad (1)$$

where $f(\tilde{\nu})$ is the spectral distribution of the fluorescence intensity, $\tilde{\nu}$ the wavenumber, $\varepsilon(\tilde{\nu})$ the extinction coefficient, and c is the velocity of light. The functions $F(\tilde{\nu})$ versus $\tilde{\nu}$ are straight lines, the slopes of which allow for the determination of T . This method gives the weighted average over the lifetime of the excited state of the luminescent molecule studied.

If the occupation density of vibrational levels in the excited state does not correspond to the temperature T of the solution studied but to a higher “vibrational temperature” T^* , (1) should be written in the modified form [5]

$$F(\tilde{\nu}, T) \equiv \ln \frac{f(\tilde{\nu}, T^*)}{\tilde{\nu}^3 \cdot \varepsilon(\tilde{\nu}, T^*)} = -\frac{hc\tilde{\nu}}{kT^*} + \text{const}. \quad (2)$$

Relation (2) has been investigated for liquid, viscous and rigid solutions [6–17]. It has been found that only in viscous and rigid solutions as well as in mixed solutions the

local temperature T^* is always higher than the ambient temperature T . Experimental investigations of (1) performed by us in the seventies [15–17] showed that both for polar and nonpolar luminescent molecules in polar and nonpolar liquid solvents the temperature T^* is equal to the ambient temperature T (within the ± 10 K accuracy). This concerns the case in which the decay time τ_F of the excited state is so long (about 10^{-9} s) that thermal equilibration occurs between the luminescent and solvent molecules.

If, however, τ_F is comparable to the relaxation time τ_r of the vibrational energy, emission occurring before relaxation is expected and the emission band is broadened on the long wavelength side. In this case the spectral distribution of emission intensity will be slightly dependent on the wavelength λ_{exc} of the excitation light [18].

Recently, Sawicki, and Knox [19, 20] introduced the term of spectral temperature $T^*(\tilde{\nu})$ defined through the local slope of $F(\tilde{\nu})$. Equation (3) can be obtained by differentiating eq. (2) versus $\tilde{\nu}$:

$$T^*(\tilde{\nu}) = -\frac{hc}{k} \left(\frac{dF(\tilde{\nu})}{d\tilde{\nu}} \right)^{-1}. \quad (3)$$

In the present work we will show that for luminescent molecules with a very short lifetime τ_F comparable to that of the relaxation time τ_r in liquid solvents of low viscosity the inequality $T^* > T$ is also valid.

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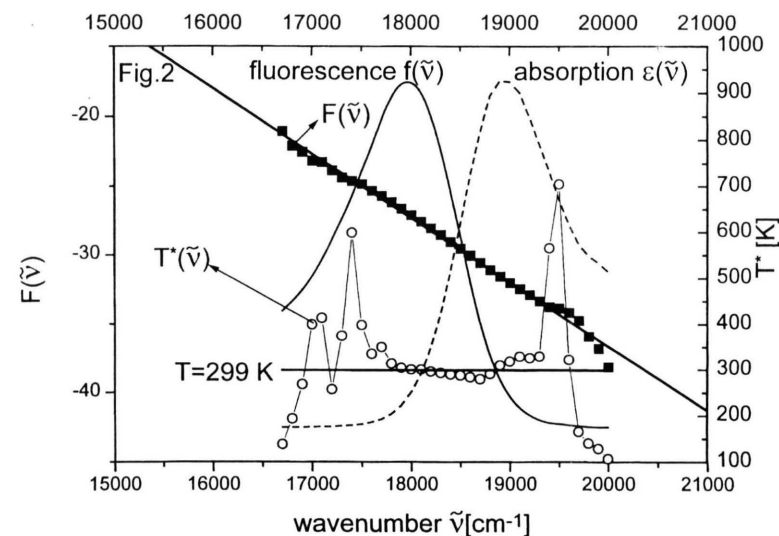
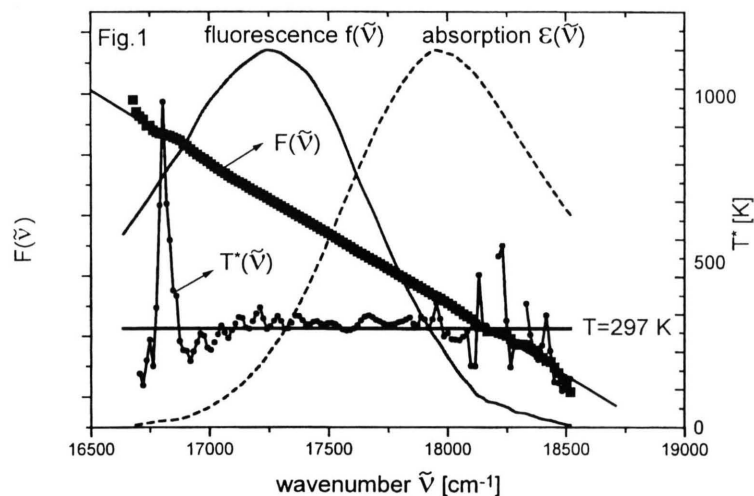


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Figs. 1 and 2. Absorption, fluorescence and $T^*(\tilde{\nu})$ spectra of rhodamine B and rhodamine 6G in water (+ a trace of 0.1 N solution of HCl), respectively, as well as the corresponding $F(\tilde{\nu})$ function).

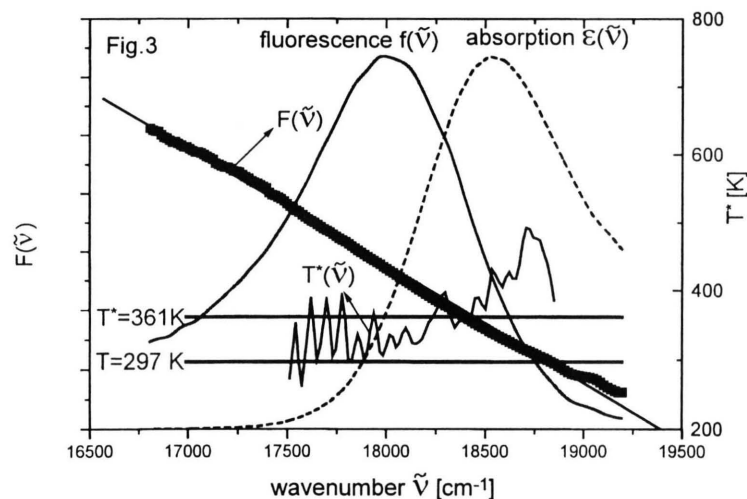


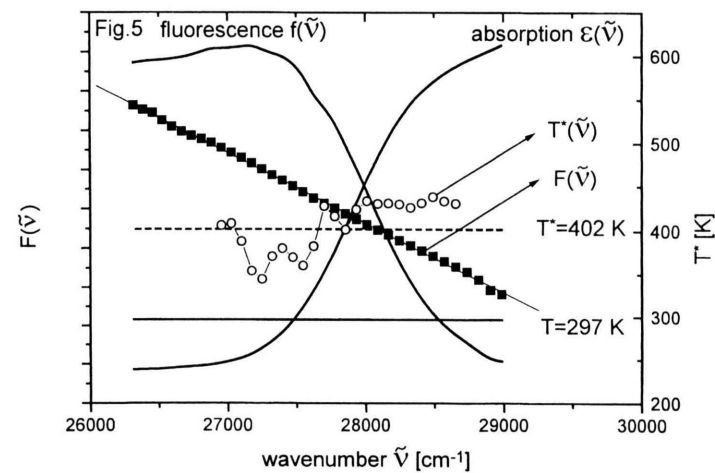
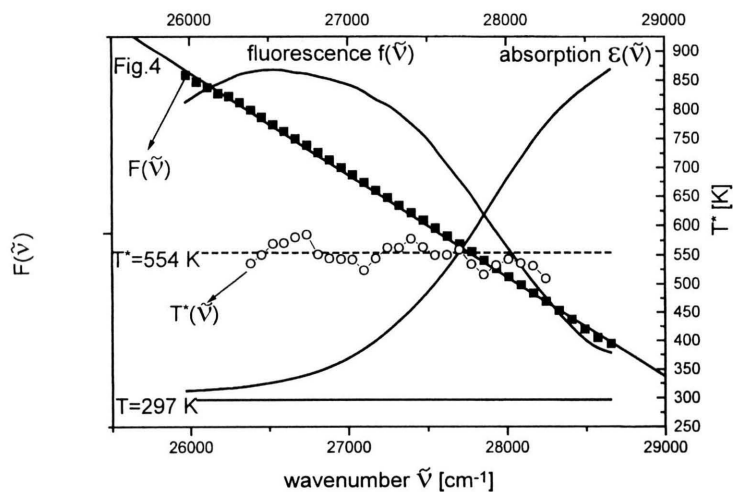
Fig. 3. Absorption, fluorescence and $T^*(\tilde{\nu})$ spectra of erythrosin B in water, as well as the corresponding $F(\tilde{\nu})$ function.

Table 1. Spectral properties and local temperatures of compounds studied in liquid solutions.

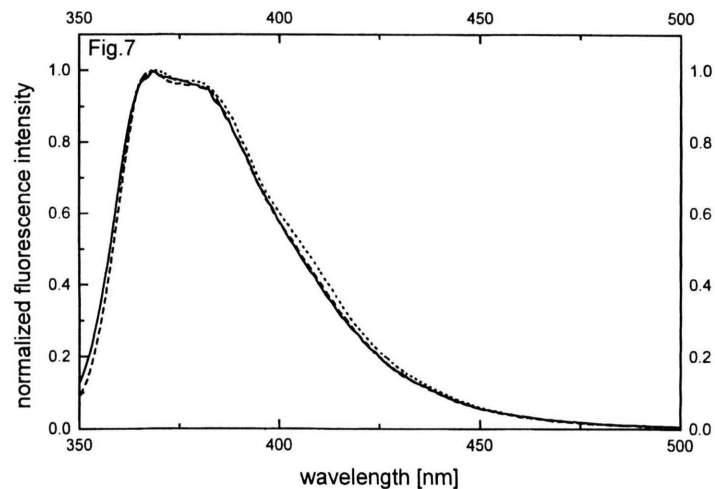
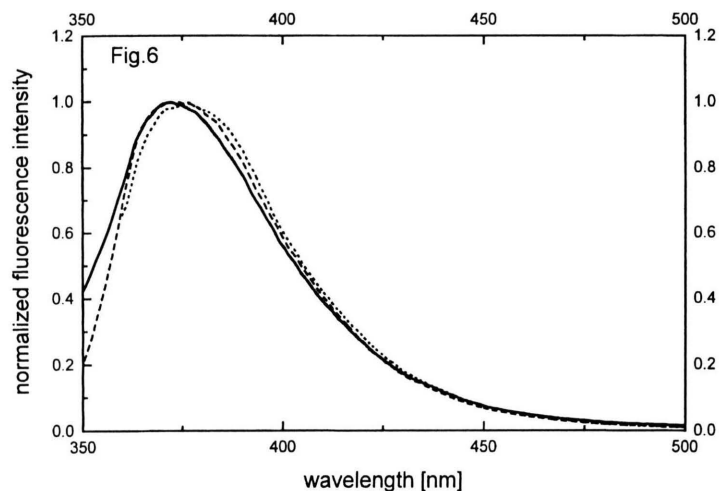
Compound	Solvent	Viscosity η (10^{-3} Pa s) ¹	τ_F [10^{-12}]	$\tilde{\nu}_{0-0}$ [cm^{-1}]	$\tilde{\nu}_{exc}$ [cm^{-1}]	T (ambient) [K]	T^* [K]	ΔT^* $= T^* - T$ [K] ²
Rhodamine B	water+HCl		1580 [23]	17614	18231	297	304	7
Rhodamine 6G	water+HCl		3950 [23]	18501	20000	299	309	10
Rhodamine S	water+HCl		?	17984	20000	298	306	8
Na-fluorescein	water+NaOH	0.89	4270	20220	22222	297	297	0
Erythrosin B	water		77 [24]	18295	19608	297	361	64
2a	<i>n</i> -hexane	0.326	7 [25]	27861	30303	297	554	257
	<i>n</i> -hexadecane	2.33	20 [25]	28108	30303	297	486	189
3a	<i>n</i> -hexane	0.326	19 [25]	27994	30303	297	402	105
	<i>n</i> -decane	0.818	26 [25]	27799	30303	297	399	102
	<i>n</i> -hexadecane	2.33	29 [25]	27710	30303	297	391	94

¹ 1 cP = 10^{-3} kg/(m s) = 10^{-3} Pa s.

² The error of the method: ± 10 K.



Figs. 4. and 5. Absorption, fluorescence and $T^*(\tilde{\nu})$ spectra of 2a and 3a in *n*-hexane, respectively, as well as the corresponding $F(\tilde{\nu})$ function.



Figs. 6. and 7. Fluorescence spectra of 2a and 3a in *n*-hexane, respectively, for different excitation wavelengths: $\lambda_{\text{exc}} = 310$ nm (solid line), $\lambda_{\text{exc}} = 330$ nm (dashed line) and $\lambda_{\text{exc}} = 350$ nm (dots).

2. Experimental

4-dimethylamino- ω -diphenylphosphinyl-trans-styrene (2a) and 4-dimethylamino- ω -methylsulphonyl-trans-styrene (3a) of spectroscopic grade were synthesized and kindly gifted by Dr. Dieter Gloyna (Humboldt Universität zu Berlin).

The dyes: rhodamine B, rhodamine S, rhodamine 6G, Na-fluorescein, and erythrosin B analytically pure were additionally purified by multiple crystallization. All alkane solvents were spectroscopically pure. Water was doubly distilled. The absorption spectra were recorded on a fully computerized Zeiss model M-40 spectrophotometer and fully corrected fluorescence spectra were measured using a spectrofluorimeter designed and built in our laboratory [21].

3. Results and Discussion

We selected seven luminescent compounds, four of them with long lifetimes (several nanoseconds) and three with short lifetimes ranging (from 7 ps to 77 ps). Figs. 1 and 2 present the absorption and fluorescence spectra as well as spectral temperature $T^*(\bar{\nu})$ and $F(\bar{\nu})$ courses for rhodamine B and rhodamine 6G, respectively, in water (a trace amount of 0.1 N solution of HCl was added to ensure the presence of a single ionic form). One can observe a linear dependence of $F(\bar{\nu})$ over the overlap region between the absorption and fluorescence band, which is consistent with (1). However, in the bands overlap a strong fluctuation of $T^*(\bar{\nu})$ around the horizontal line corresponding to the temperature of the experiment. Strong departures of this function from linearity appear only on the edges of the absorption and fluorescence bands (for example for rhodamine B above $\bar{\nu} = 18400 \text{ cm}^{-1}$ (fluorescence band) and below $\bar{\nu} = 17100 \text{ cm}^{-1}$ (absorption band)). A similar behaviour is found for rhodamine S and Na-fluorescein in water. Temperatures T^* determined based on (2) are listed in Table 1. It can be seen that the differences $\Delta T = T^* - T$ are within the limits of experimental error $\pm 10 \text{ K}$ (resulting from the uncertain-

ty of the slope of $F(\bar{\nu})$). Our results are in a full agreement with the earlier ones obtained for nonpolar and polar luminescent molecules in nonpolar and polar solvents [15–17, 22]. It has been found therein that the polarity of both the luminescent and solvent molecules does not exhibit a significant dependence on the local temperature.

A completely different behaviour, compared to the mentioned dyes, was found in this work for other luminescent compounds, the lifetimes of which are very short. Figures 3–5 show the results of investigations on erythrosin B in water and on 2a and 3a in *n*-hexane, respectively. The linear dependence of $F(\bar{\nu})$ according to (2) is fulfilled over the absorption and fluorescence overlap range, however the temperatures T^* determined from the slope of $F(\bar{\nu})$ exceed distinctly the ambient temperature T (cp. Table 1). In Figs. 3–5 the function $T^*(\bar{\nu})$ is plotted only over such a range of absorption and fluorescence where no significant experimental error is made.

Data obtained for these short lived solutes suggest that higher values of T^* are caused by the non-establishment of statistical equilibrium before the emission act over the vibrational levels of excited molecules. In this case part of the excess of excitation energy is not dispersed throughout the solvent during the short lifetime of the solute. This is the reason of the high local temperature T^* which holds also during the emission process. In view of the very short lifetimes of the molecules studied (7 ps–77 ps) and the lack of thermal equilibrium with the surrounding medium it should be expected that the intensity distribution in the emission band as well as its spectral location will not be independent of the excitation light frequency. Figures 6 and 7 show that indeed a slight dependence of the fluorescence band on the excitation light frequency is observed.

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